

Phase diagrams of Ising mixtures



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Mean Field Theory¹

We describe a mixture of a van der Waals fluid and a ferromagnetic Ising fluid with mole fraction x and magnetization m at zero magnetic field in the framework of mean field theory. Setting $x = 1$ yields the mean field theory of the pure Ising fluid.

- Equations of state:

$$p_r = \frac{T_r}{V_r - 1} - 54 \alpha(x, m) \quad m = \tanh\left(\frac{27R_m xm}{4 V_r T_r}\right)$$

- Quadratic mixing rule:

$$\alpha(x, m) = \left(1 + R_m m^2\right) x^2 + 2 \frac{1 - \Lambda}{1 + \zeta} x (1 - x) + \frac{1 - \zeta}{1 + \zeta} (1 - x)^2$$

- 3 parameters characterizing the mixture:

$$\zeta = \frac{a_{22} - a_{11}}{a_{11} + a_{22}}, \quad \Lambda = \frac{a_{11} - 2a_{12} + a_{22}}{a_{11} + a_{22}}, \quad R_m = \frac{1}{2} \frac{a_m}{a_{22}}$$

MC Simulation of phase transitions²

Potential between particles i and j of species a and b : Lennard-Jones + anisotropic Yukawa

$$u_{ij}(r) = \begin{cases} 4\epsilon_{ab} \left[\left(\frac{\sigma}{r}\right)^{12} - \left(\frac{\sigma}{r}\right)^6 \right] - R\varepsilon_{22} s_i s_j e^{-r/\sigma} & r < r_c \\ 0 & r > r_c \end{cases} + \text{tail correction}$$

with spins $s_i = \pm 1$ for species 2, $s_i = 0$ for species 1

→ obtain 2D-histograms $H_i(S_1, S_2)$ from i -th simulation of length n_i at K_1^i, K_2^i

→ Multiple histogram reweighting:

$$P_{K_1 K_2}(S_1, S_2) = \frac{\sum_{i=1}^n H_i(S_1, S_2)}{\sum_{j=1}^n n_j e^{(K_1^j - K_1)S_1 + (K_2^j - K_2)S_2 - C_j}}, \quad e^{C_j} = \sum_{S_1, S_2} P_{K_1^j K_2^j}(S_1, S_2)$$

Pure fluid

Grand Canonical Ensemble

- constant T, V, μ
- box size $L = 10\sigma \dots 16\sigma$
- find shape of binodal near criticality, extrapolate tricritical point

$$S_1 = E \quad K_1^j = -\beta_j \quad S_2 = N \quad K_2^j = \beta_j \mu_j$$

Density of States MC³

- for high density/low temperature regions in pure fluid diagrams

Canonical Ensemble

- constant N, T, V
- $N = 150, 300, 500$, finite size scaling
→ find para-ferro phase transitions

$$S_1 = E \quad K_1^j = -\beta_j \quad S_2 = M \quad K_2^j = \beta_j H = 0$$

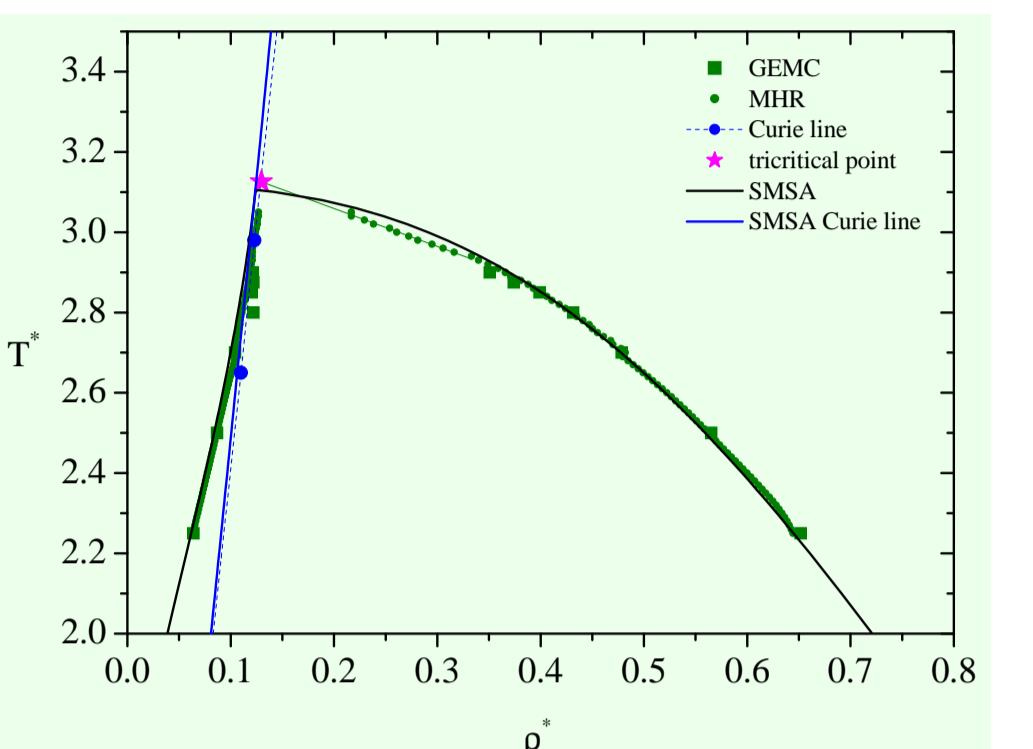
Gibbs Ensemble

- constant N, T, V, p
- $N = 500 \dots 1000$
- not suitable for low temperatures
- does not work near criticality

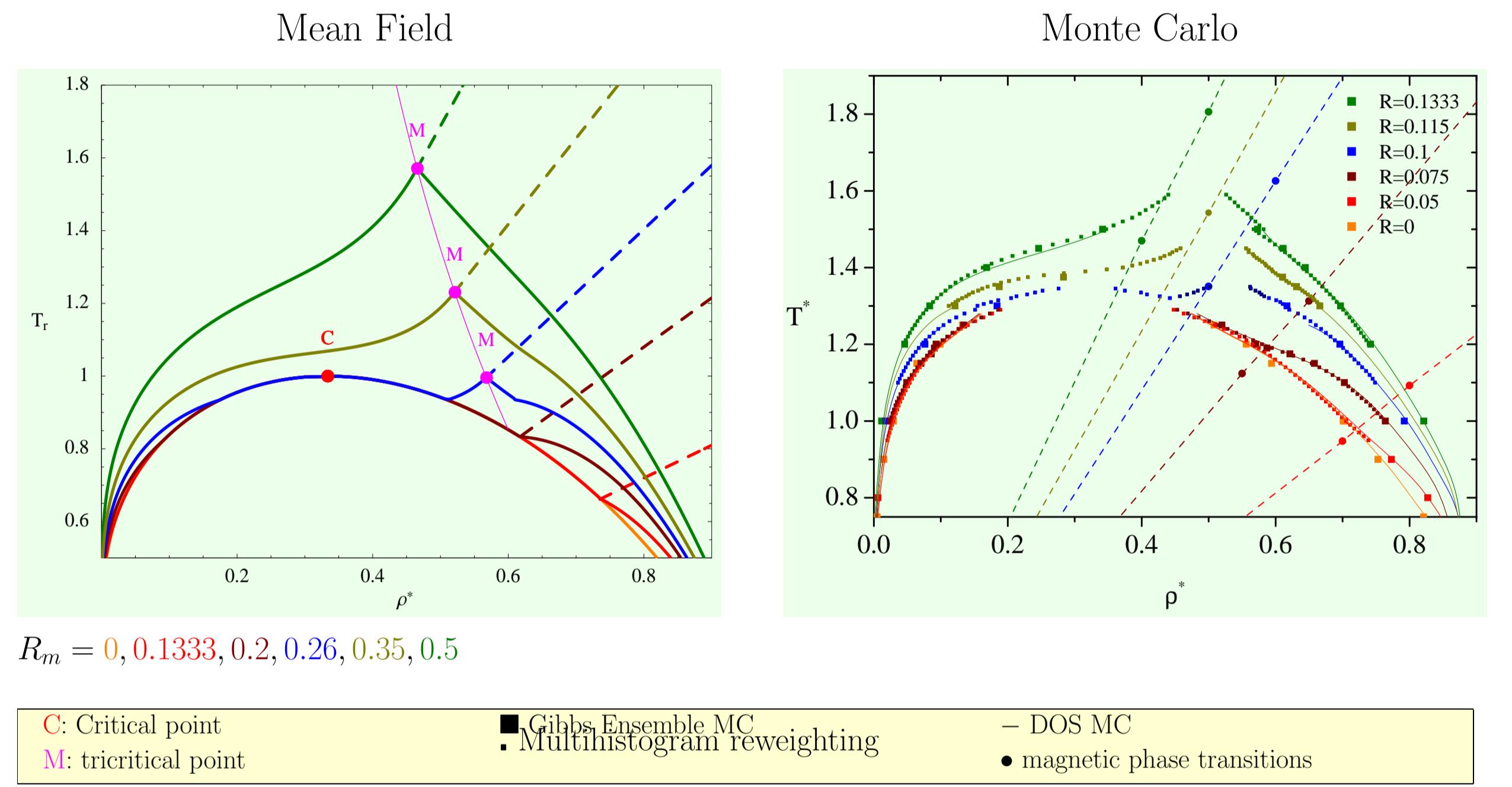
Results for pure Ising fluids

ideal Ising fluid:

- no nonmagnetic attraction: $\varepsilon_{22} = 0, R \rightarrow \infty, \varepsilon_m \equiv R\varepsilon_{22} = \text{const.}$
- soft-core repulsion: $\varphi_{SC}(r) = \begin{cases} 4\varepsilon_{sc} \left[\left(\frac{\sigma}{r}\right)^{12} - \left(\frac{\sigma}{r}\right)^6 \right] + \varepsilon_{sc} & r < \sqrt[6]{2}\sigma \\ 0 & r > \sqrt[6]{2}\sigma \end{cases}$
- comparison with integral equation theory:
Soft Mean Spherical Approximation (SMSA)⁴

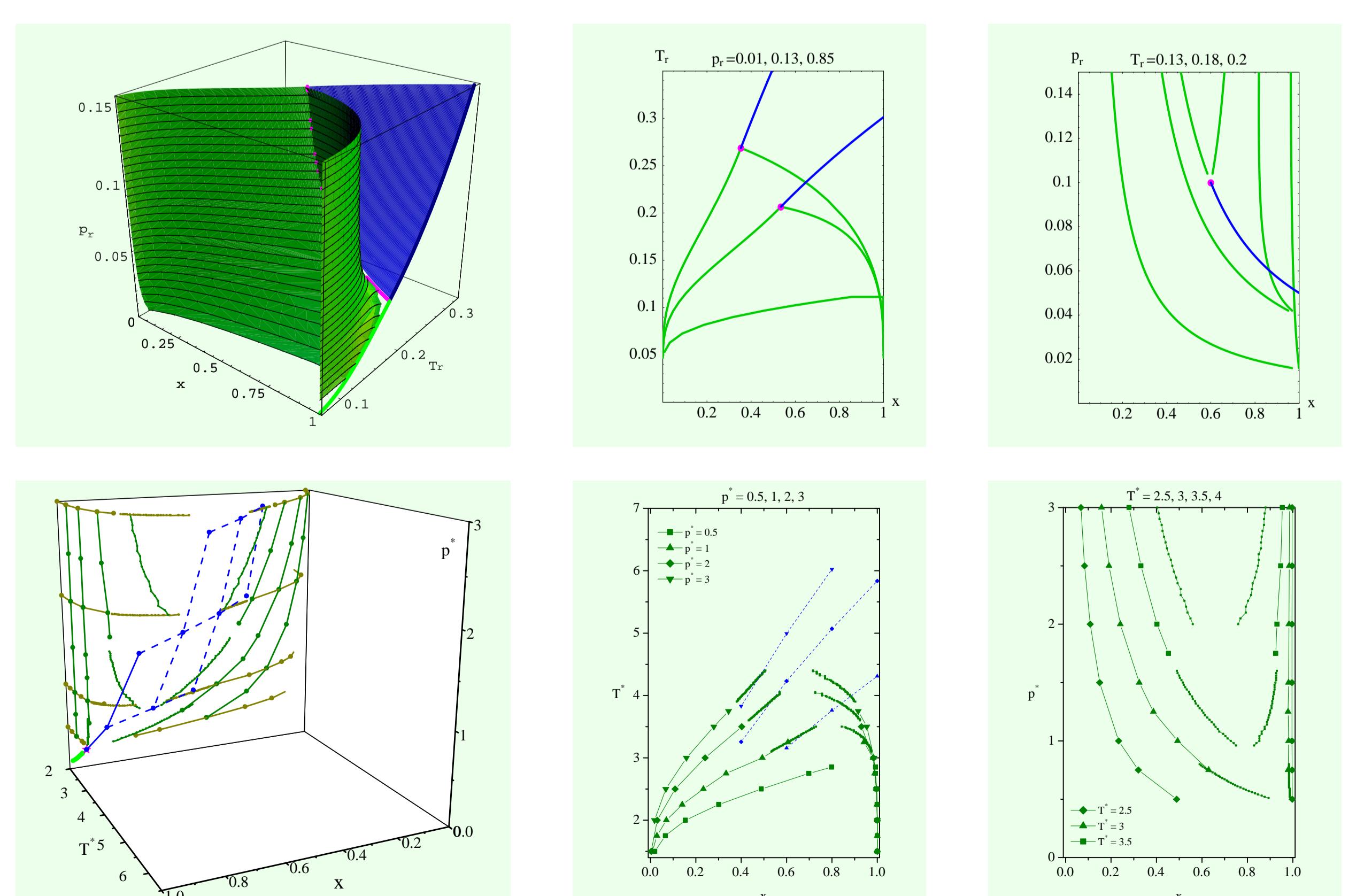


van der Waals Ising fluid:



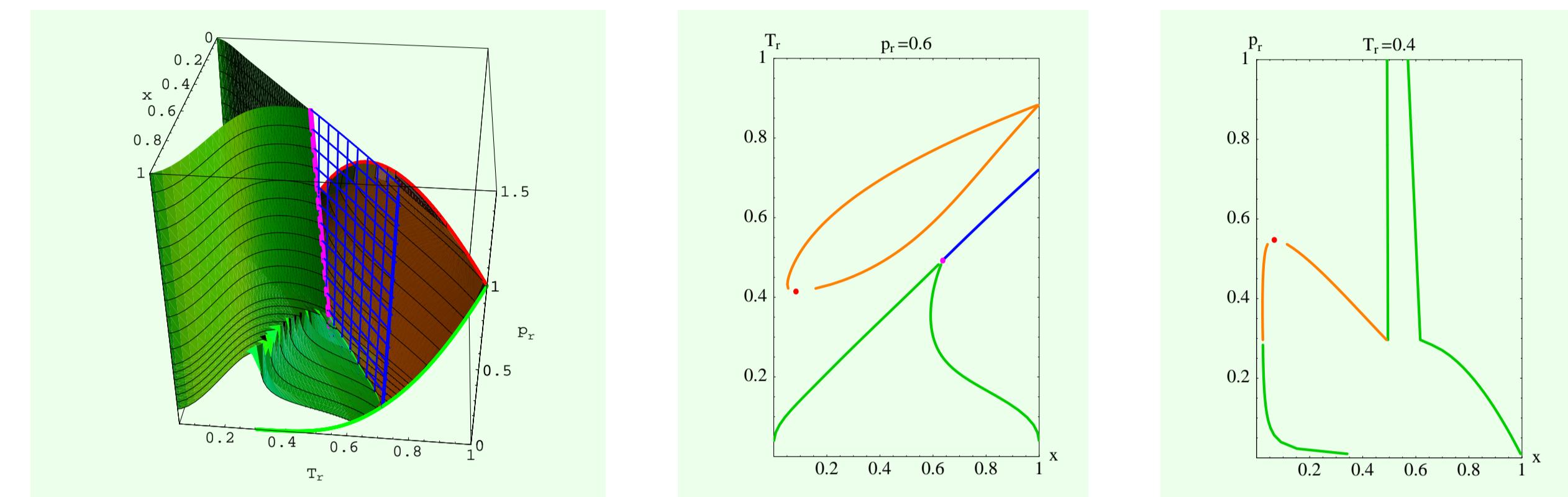
Results for the ideal Ising mixture

ideal Ising fluid + SC fluid

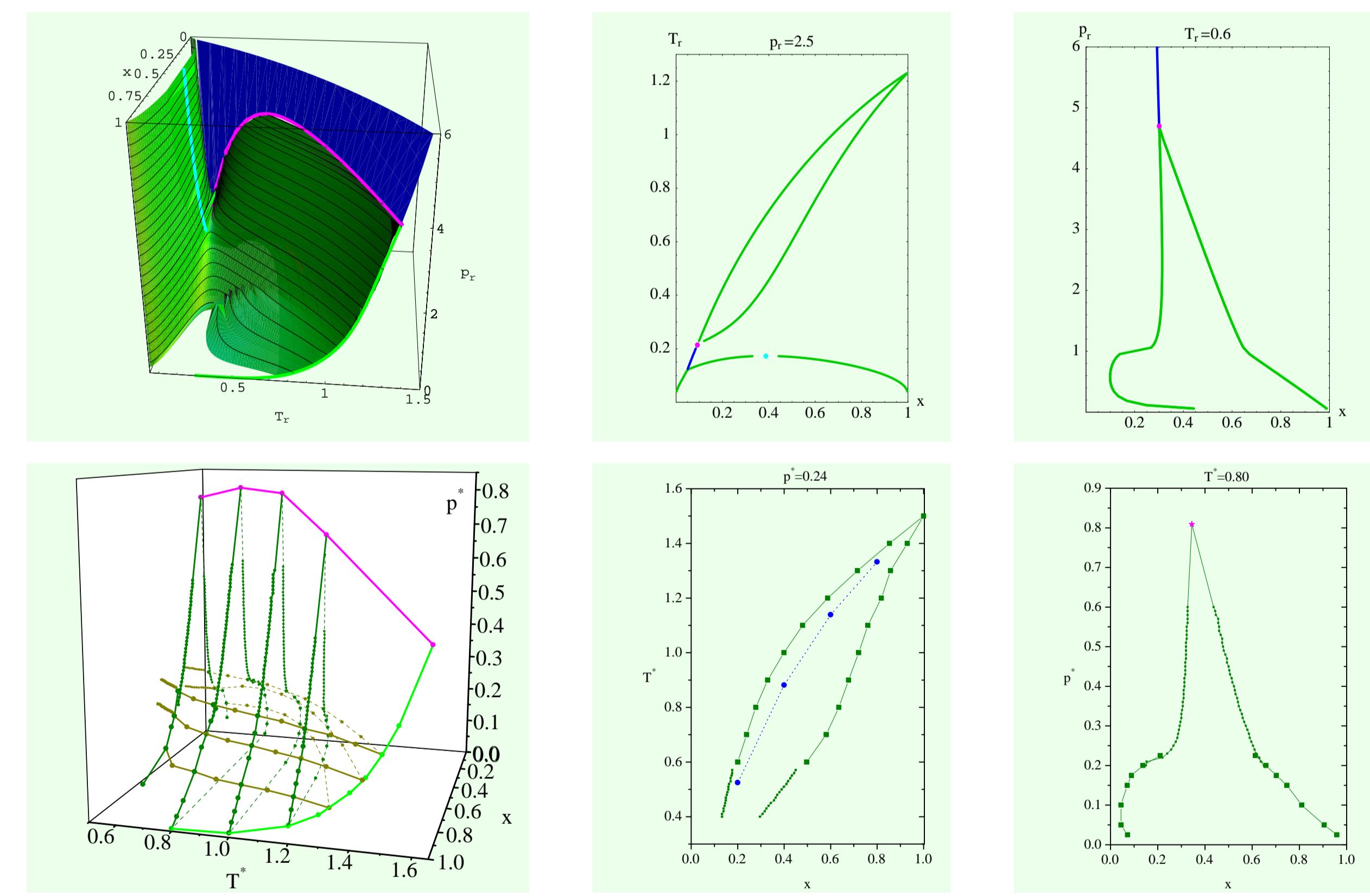


Results for Ising mixtures

$\zeta = 0.5, \Lambda = -0.05, R_m = 0.15$



$\zeta = 0.5, \Lambda = -0.25, R_m = 0.4$



Summary

While in ordinary binary fluid mixtures tricritical points occur only under special circumstances, mixtures with a magnetic fluid component show a line of tricritical points, that has either the character of a consolute point line or a plait point line. These topologies of the phase diagrams calculated in mean field theory were verified with Monte Carlo simulations at certain parameter values. (Grant: FWF P15247)

References

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- [2] W. Fenz, R. Folk, Cond. Matt. Phys. **6**, 675 (2003)
- [3] F. G. Wang and D. P. Landau, Phys. Rev. Lett. **86**, 2050 (2001)
- [4] I. P. Omelyan, I. M. Mryglod, R. Folk, W. Fenz, submitted to Phys. Rev. E