

21. The Schrödinger equation for periodic potentials I

- (a) Show that for a periodic potential with period  $a$ ,  $V(x + a) = V(x)$ , the Hamiltonian commutes with the translation operator  $T_a$ ,  $[T_a, H] = 0$ . The operator  $T_a$  translates the whole system by a distance  $a$  and is defined by

$$T_a\psi(x) = \psi(x + a).$$

What is the consequence of this commutator relation?

- (b) Show that from part (a) it follows that the solutions  $\psi(x)$  of the Schrödinger equation for a periodic potential must be of the form

$$\psi(x) = e^{ikx}u(x) \quad \text{with} \quad u(x) = u(x + h) \quad (1)$$

and a real-valued parameter  $k$ . This relation is called the *Bloch theorem*.

*Hint:* Use the fact that the translation operator does not change the norm of the wavefunction. What do you know about the operator product  $T_aT_{a'}$ ?

- (c) Consider a periodic system where the potential is so weak that you can neglect it — *i.e.* you can assume  $V(x) = 0$ , but the wavefunctions are still constrained to have the form given by equation (1). Calculate the eigenenergies of the system in dependence of the parameter  $k$ . Make a sketch of  $E(k)$  in the region  $[-\frac{\pi}{a}, +\frac{\pi}{a}]$ . Why is it sufficient to constrain  $k$  to this region?

*Hint:* As  $u(x + a) = u(x)$ , the function  $u(x)$  can be expressed as a Fourier series,

$$u(x) = \sum_{n=-\infty}^{\infty} u_n e^{iG_n x}, \quad \text{with} \quad G_n = \frac{2\pi}{a}n.$$

22. The Schrödinger equation for periodic potentials II

Now consider a one-dimensional quantum system in the periodic potential

$$V(x) = -V_0 \sum_{n=-\infty}^{+\infty} \delta(x - na), \quad V_0 > 0.$$

Derive an equation for the energy eigenvalues of the problem, and show how the equation can be solved graphically. Are there solutions for any value of  $k$ ? Try to make a sketch of the function  $E(k)$ .

*Hint:* Use the usual Ansatz for the wavefunction  $\psi(x)$  in the interval  $[0, a]$ . Then use the following conditions to determine the unknowns:

- Continuity of  $\psi(x)$  at  $x = 0$ .
- The jump in  $\psi'(x)$  at  $x = 0$  (see problem 14 (b)).
- The Bloch theorem (equation (1) above).

23. A basis of angular momentum eigenstates is given by  $\{|\tau l m_z\rangle\}$  ( $l, m_z$  quantum numbers of  $L^2$  and  $L_z$ ;  $\tau$  is a set of quantum numbers, which belong to a set of self-adjoint operators  $T$  commuting with all components of  $\vec{L}$ .  $\tau, l, m$  form a complete set of quantum numbers).
- (a) In which of the states  $|\tau l m_z\rangle$  with fixed  $l$  does the uncertainty of the components  $L_x$  and  $L_y$  have the smallest value? Calculate this value.
- (b) Are there states in the considered state space, where all components of the concerning angular momentum have a 'sharp value' (*i.e.* no uncertainty)?
24. Consider an arbitrary physical system, whose four-dimensional state space is spanned by the four eigenstates  $|l, m_z\rangle$  of the operators  $L^2$  and  $L_z$  ( $l = 0, 1; -l < m_z < l$ ). The eigenvalues are  $l(l+1)\hbar^2$  and  $m_z\hbar$ , respectively, so that

$$\begin{aligned} L_{\pm}|l, m_z\rangle &= \hbar\sqrt{l(l+1) - m_z(m_z \pm 1)}|l, m_z \pm 1\rangle \\ L_+|l, l\rangle &= L_-|l, -l\rangle = 0. \end{aligned} \quad (2)$$

- (a) Write the eigenstates  $|l, m_x\rangle$  of  $L^2$  and  $L_x$  in terms of states  $|l, m_z\rangle$ .
- (b) Consider a system in the normalized state (notation  $|l, m_z\rangle$ )

$$|\Psi\rangle = \alpha|1, 1\rangle + \beta|1, 0\rangle + \gamma|1, -1\rangle + \delta|0, 0\rangle.$$

What are the probabilities, that measurements of  $L^2$  and  $L_z$  give  $2\hbar^2$  and  $\hbar$ , respectively? Calculate also the average value of  $L_z$ , when the system is in the state  $|\Psi\rangle$ , and the probabilities of each possible value this observable can have.